

# Transient Phenomena in Electric Circuits

## Growth and Decay of Currents in series LR Circuits

(i) Growth of current

induction of back emf  $-L \frac{di}{dt}$   
opposes the driving emf  $V$

$$i = \frac{V - L \frac{di}{dt}}{R} \quad \text{or} \quad L \frac{di}{dt} + Ri = V$$

$$\text{or} \quad \frac{di}{i - V/R} = -\frac{R}{L} dt$$

Assuming that  $i = 0$  at  $t = 0$

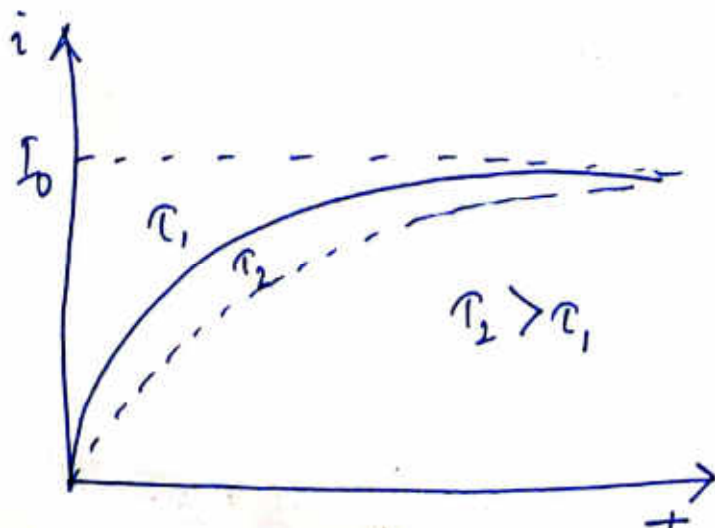
$$\int_0^i \frac{di}{i - V/R} = -\frac{R}{L} \int_0^t dt$$

$$\text{or} \quad \ln \frac{i - V/R}{-V/R} = -\frac{R}{L} t$$

$$\text{or} \quad i = \frac{V}{R} \left( 1 - e^{-R/L t} \right)$$

$$= I_0 \left( 1 - e^{-R/L t} \right)$$

$$\frac{di}{dt} = I_0 \frac{R}{L} e^{-R/L t} = \frac{R}{L} (I_0 - i)$$



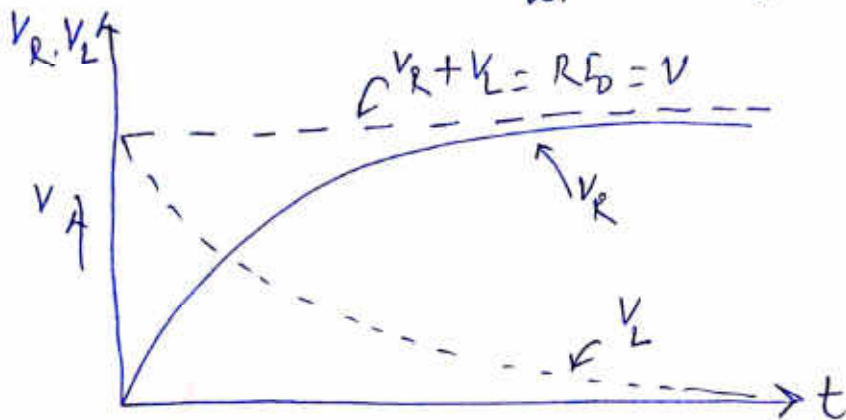
At  $t = \tau$ ,

$$i = I_0 \left( 1 - \frac{1}{e} \right) \approx 0.63 I_0$$

When the current reaches its max<sup>m</sup> steady value the variation of current with time stops and so the back emf becomes zero. At this stage the current behaves as if there were no inductor in the circuit. The instantaneous voltage drops across L and R are given by

$$V_R = Ri = R I_0 (1 - e^{-R/Lt})$$

$$\text{and } V_L = L \frac{di}{dt} = R I_0 e^{-R/Lt}$$



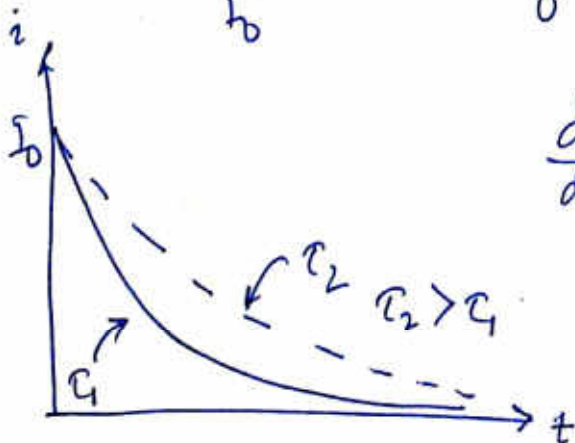
(i) Decay of current

$$i = \frac{0 - L \frac{di}{dt}}{R} \quad \text{or} \quad L \frac{di}{dt} + Ri = 0 \quad \text{or} \quad \frac{di}{i} = -\frac{R}{L} dt$$

Assuming that  $i = I_0$  at  $t = 0$

$$\int_{I_0}^i \frac{di}{i} = -\frac{R}{L} \int_0^t dt \quad \text{or} \quad \ln \frac{i}{I_0} = -\frac{R}{L} t$$

$$\text{or} \quad i = I_0 e^{-R/Lt}$$



$$\frac{di}{dt} = -\frac{R}{L} I_0 e^{-R/Lt} = -\frac{I_0}{\tau} e^{-t/\tau}$$

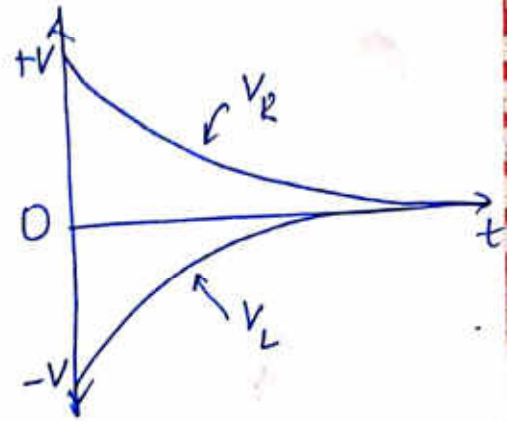
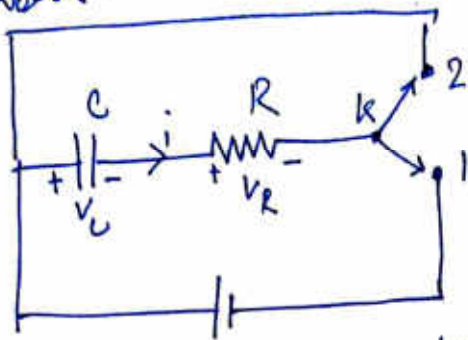
$$\text{or} \quad \left. \frac{di}{dt} \right|_{t=0} = -\frac{I_0}{\tau}$$

$$\text{at } t = \tau, i = I_0/e \approx 0.37 I_0.$$

$$V_L = L \frac{di}{dt} = -R I_0 e^{-R/L t} = -V e^{-R/L t}$$

$$V_R = R i = R I_0 e^{-R/L t} = V e^{-R/L t}$$

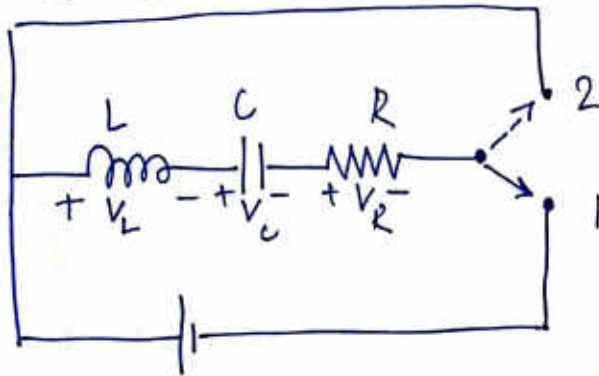
Homework:



$$q = q_0 (1 - e^{-t/CR})$$

$$q = q_0 e^{-t/CR}$$

Series LRC Circuit



$$L \frac{di}{dt} + \frac{q}{C} + Ri = V$$

$$i = \frac{dq}{dt}$$

$$L \frac{d^2q}{dt^2} + R \frac{dq}{dt} + \frac{1}{C} q = V$$

$$\omega \frac{d^2q}{dt^2} + 2b \frac{dq}{dt} + \omega_0^2 \left( q - \frac{V}{\omega_0^2} \right) = 0$$

$$2b = \frac{R}{L}$$

$$\omega_0^2 = \frac{1}{LC}$$

$$\Rightarrow \frac{d^2Q}{dt^2} + 2b \frac{dQ}{dt} + \omega_0^2 Q = 0$$

$$Q = q - \frac{V}{\omega_0^2}$$

Trial sol:  $Q = A e^{\alpha t}$

$$(\alpha^2 + 2b\alpha + \omega_0^2) A e^{\alpha t} = 0$$

$$\alpha^2 + 2b\alpha + \omega_0^2 = 0$$

$$\alpha_1 = -b + \sqrt{b^2 - \omega_0^2} \quad \text{and} \quad \alpha_2 = -b - \sqrt{b^2 - \omega_0^2}$$

$$Q = A_1 e^{\alpha_1 t} + A_2 e^{\alpha_2 t}$$

$$= e^{-bt} \left[ A_1 e^{\sqrt{b^2 - \omega_0^2} t} + A_2 e^{-\sqrt{b^2 - \omega_0^2} t} \right]$$

$$q = Q + \frac{V}{L\omega_0^2} = q_0 + e^{-bt} [A_1 e^{pt} + A_2 e^{-pt}]$$

$$q_0 = V/L\omega_0^2 = V.C \text{ and } p = \sqrt{b^2 - \omega_0^2}$$

$$t=0, q=0, i = \frac{dq}{dt} = 0$$

$$0 = q_0 + A_1 + A_2 \text{ and } 0 = (A_1 - A_2)p - (A_1 + A_2)b$$

$$A_1 = -\frac{q_0}{2} \left(1 + \frac{b}{p}\right)$$

$$A_2 = -\frac{q_0}{2} \left(1 - \frac{b}{p}\right)$$

$$\text{Complete sol}^n: q = q_0 - \frac{q_0}{2} e^{-bt} \left[ \left(1 + \frac{b}{p}\right) e^{pt} + \left(1 - \frac{b}{p}\right) e^{-pt} \right]$$

$$\text{Case I: } b^2 > \omega_0^2$$

$$p = \sqrt{b^2 - \omega_0^2}$$

$$\text{Case II: } b^2 = \omega_0^2$$

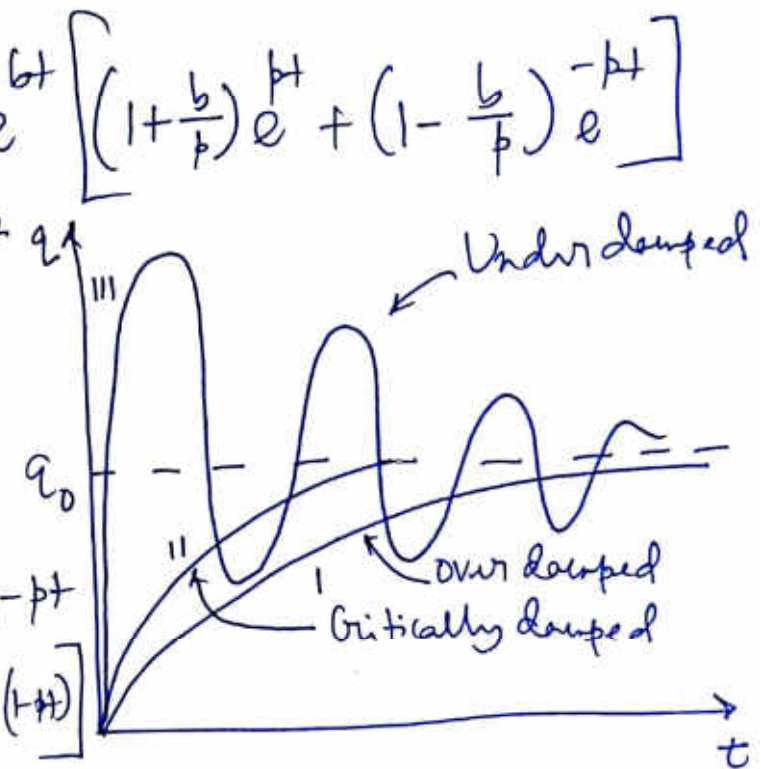
$b > 0$ ,  $b/p$  tends to  $\infty$

we assume  $p \ll 1$

$$e^{pt} \approx 1 + pt \text{ and } e^{-pt} = 1 - pt$$

$$q = q_0 - \frac{q_0}{2} e^{-bt} \left[ \left(1 + \frac{b}{p}\right)(1 + pt) + \left(1 - \frac{b}{p}\right)(1 - pt) \right]$$

$$= q_0 - q_0 (1 + bt) e^{-bt}$$



$b^2 = \omega_0^2$  or  $R = 2\sqrt{L/C}$  Critical damping resistance

$$\text{Case III: } b^2 < \omega_0^2$$

$$p = \sqrt{b^2 - \omega_0^2} \text{ is imaginary}$$

$$p = j\sqrt{\omega_0^2 - b^2} = j\omega, \omega = \sqrt{\omega_0^2 - b^2}$$

$$q = q_0 - \frac{q_0}{2} e^{-bt} \left[ \frac{e^{j\omega t} - e^{-j\omega t}}{e + e} + \frac{b}{j\omega} \left( \frac{e^{j\omega t} - e^{-j\omega t}}{e - e} \right) \right]$$

$$= q_0 - q_0 e^{-bt} \left[ \cos \omega t + \frac{b}{\omega} \sin \omega t \right]$$

$$= q_0 - \frac{q_0}{\omega} e^{-bt} \left[ \omega \cos \omega t + b \sin \omega t \right]$$

putting  $\omega = A \cos \phi$ ,  $b = A \sin \phi$

$$A = \sqrt{\omega^2 + b^2} = \omega_0 \text{ and } \phi = \tan^{-1} \frac{b}{\omega}$$

$$q = q_0 - q_0 \frac{\omega_0}{\omega} e^{-bt} \cos(\omega t - \phi) \quad (4)$$

freq<sup>y</sup> of  $Qe^{bt}$

$$f = \frac{\omega}{2\pi} = \frac{\sqrt{\omega_0^2 - b^2}}{2\pi} = \frac{1}{2\pi} \sqrt{\frac{1}{LC} - \frac{R^2}{4L^2}}$$

$$i = \frac{dq}{dt} = \frac{q_0}{\omega} b e^{-bt} [\omega \cos \omega t + b \sin \omega t] - \frac{q_0}{\omega} e^{-bt} [-\omega^2 \sin \omega t + b \omega \cos \omega t]$$

$$= \frac{q_0 \omega^2}{\omega} e^{-bt} \sin \omega t$$

R becomes negligible then  $b=0$ ,  $\omega = \omega_0$

$$q = q_0 - q_0 e^{-\omega_0 t} \quad \omega_0 \rightarrow \text{undamped natural freq<sup>y</sup>}$$

$$i = q \omega_0 \sin \omega_0 t$$

Homework: Discharging of the Capacitor

Prob: Examine whether the discharge of a charged capacitor 0.1  $\mu\text{F}$  through an inductive coil of inductance 100mH and resistance 200 $\Omega$  is oscillatory or not. If it is osc<sup>ty</sup>, find the freq<sup>y</sup> of  $Qe^{bt}$ .

$$R = 2\sqrt{\frac{L}{C}} = 2\sqrt{\frac{100 \times 10^{-3}}{0.1 \times 10^{-6}}} = 2000 \Omega$$

As the coil resistance is smaller than this value the discharge will be osc<sup>ty</sup>. The freq<sup>y</sup> of  $Qe^{bt}$

$$f = \frac{1}{2\pi} \sqrt{\frac{1}{LC} - \frac{R^2}{4L^2}} = \frac{1}{2\pi} \sqrt{\frac{1}{0.1 \times 0.1 \times 10^{-6}} - \frac{200^2}{4 \times (0.1)^2}}$$

$$= 1584 \text{ Hz.}$$

Topic : Transient Phenomena in Electric Circuits

(ii) Discharging of the capacitor

Source is cut off and the stored charge on the capacitor begins to discharge through L & R.

$$L \frac{di}{dt} + \frac{q}{C} + Ri = 0$$

$$\text{or } L \frac{d^2q}{dt^2} + R \frac{dq}{dt} + \frac{1}{C}q = 0$$

$$\text{or } \frac{d^2q}{dt^2} + 2b \frac{dq}{dt} + \omega_0^2 q = 0$$

$$2b = R/L$$

$$\omega_0^2 = 1/LC$$

$$q = e^{-bt} \left[ A_1 e^{pt} + A_2 e^{-pt} \right] \quad \phi = \sqrt{b^2 - \omega_0^2}$$

$$q = \frac{q_0}{2} e^{-bt} \left[ \left(1 + \frac{b}{p}\right) e^{pt} + \left(1 - \frac{b}{p}\right) e^{-pt} \right]$$

$$A_1 = \frac{q_0}{2} \left(1 + \frac{b}{p}\right)$$

$$A_2 = \frac{q_0}{2} \left(1 - \frac{b}{p}\right)$$

Case I:  $p$  : +ve non oscillatory discharge of the capacitor

Case II:  $b^2 = \omega_0^2$

$p = 0$ ,  $b/p$  tends to  $\infty$

$p \ll 1$ ,  $e^{\pm pt} \approx 1 \pm pt$

$$q = q_0 (1 + bt) e^{-bt}$$

Case III:  $b^2 < \omega_0^2$

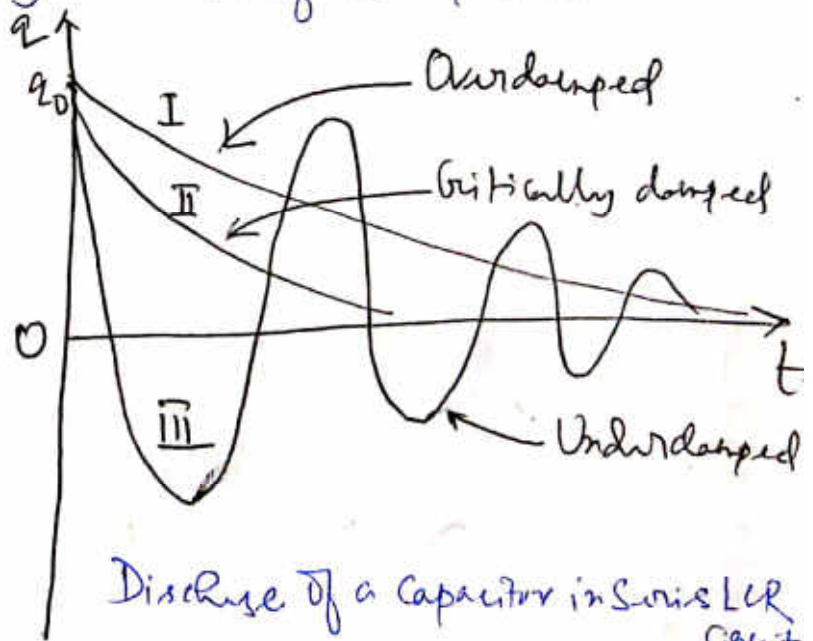
$p = \sqrt{b^2 - \omega_0^2}$  is imaginary

$$p = j\sqrt{\omega_0^2 - b^2} = j\omega, \quad \omega = \sqrt{\omega_0^2 - b^2}$$

$$q = \frac{q_0 \omega_0}{\omega} e^{-bt} \cos(\omega t - \theta), \quad \theta = \tan^{-1} \frac{b}{\omega} \quad \text{Oscillatory!}$$

$$f = \frac{\omega}{2\pi} = \frac{\sqrt{\omega_0^2 - b^2}}{2\pi} = \frac{1}{2\pi} \sqrt{\frac{1}{LC} - \frac{R^2}{4L^2}}$$

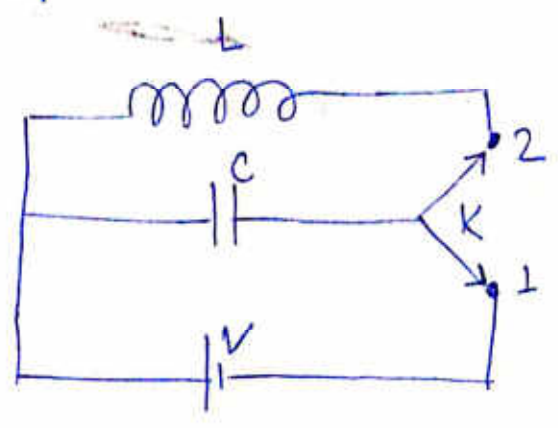
$$i = \frac{dq}{dt} = -\frac{q_0 \omega_0^2}{\omega} e^{-bt} \sin \omega t$$



Discharge of a Capacitor in Series LCR Circuit

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# Discharge of a Capacitor Through a Pure Inductor and Electrical Oscillations



$$\omega_0^2 = \frac{1}{LC}$$

$$L \frac{di}{dt} + \frac{q}{C} = 0 \quad \Rightarrow \quad L \frac{d^2q}{dt^2} + \frac{q}{C} = 0 \quad \Rightarrow \quad \frac{d^2q}{dt^2} + \omega_0^2 q = 0$$

$$q = A_1 \cos \omega_0 t + A_2 \sin \omega_0 t$$

at  $t=0$ ,  $q = q_0 = VC$  and  $\frac{dq}{dt} = 0$ ;  $q_0 = A_1$  &  $A_2 = 0$

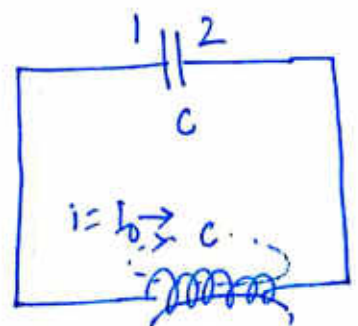
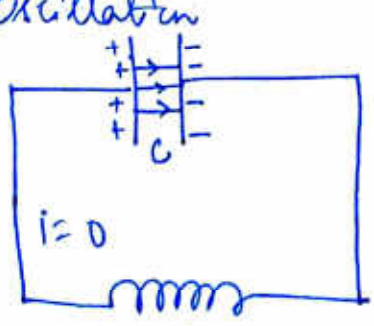
$$q = q_0 \cos \omega_0 t$$

Thus the charge on the capacitor varies in an oscillatory manner with cert amplitude  $q_0$  and  $f = \frac{\omega_0}{2\pi} = \frac{1}{2\pi\sqrt{LC}}$

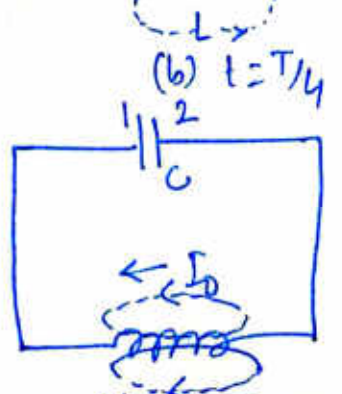
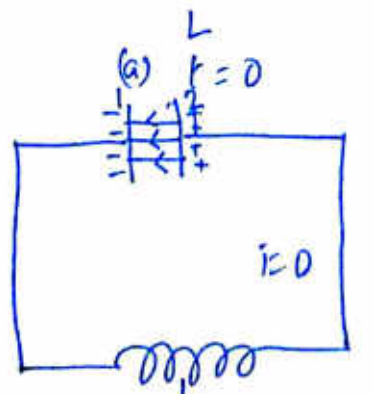
$$i = \frac{dq}{dt} = -q_0 \omega_0 \sin \omega_0 t$$

$$\Rightarrow i = q_0 \omega_0 \cos(\omega_0 t + \pi/2)$$

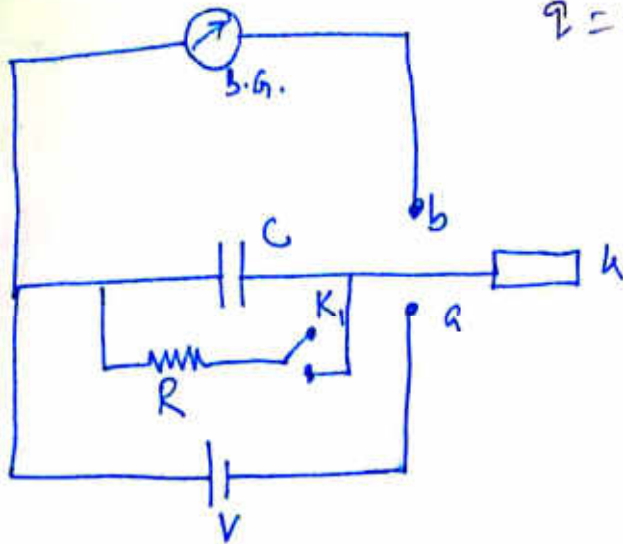
## Electric Oscillation



electrostatic en  $\frac{q_0^2}{2C}$   
magnetic en  $\frac{1}{2} L i^2$



# Measurement of High Resistance by Leakage



$$q = q_0 e^{-t/CR} \quad \therefore R = \frac{t}{C \ln \frac{q_0}{q}}$$

first ballistic throw is noted

$$\frac{q_0}{q} = \frac{d_0}{d}$$

$$R = \frac{t}{C \ln \frac{d_0}{d}}$$

$$R_0 = \frac{t}{C \ln \frac{d_0}{d}}$$

$$\frac{1}{R} = \frac{1}{R_a} + \frac{1}{R_0}$$

Natural leakage resistance

Prob: at any time t is

$$i = i_1 + i_2$$

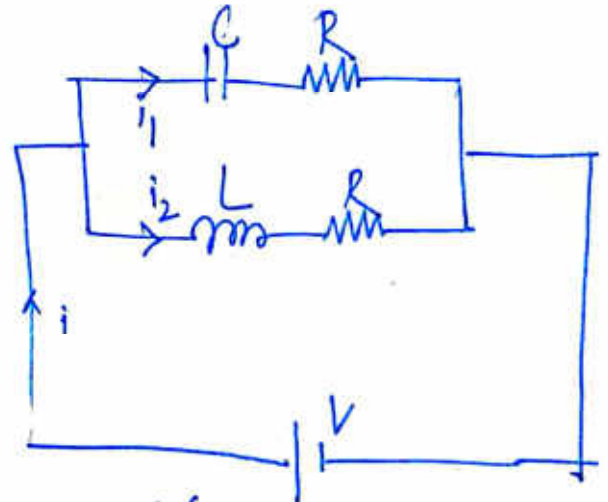
$$q = q_0 (1 - e^{-t/CR}) \quad \text{where } q_0 = CV$$

$$\therefore i_1 = \frac{dq}{dt} = \frac{q_0}{CR} e^{-t/CR} = \frac{V}{R} e^{-t/CR}$$

$$\text{also } i_2 = \frac{V}{R} (1 - e^{-R/Lt}) = \frac{V}{R} (1 - e^{-t/CR})$$

$$\text{Therefore } i = i_1 + i_2 = \frac{V}{R}$$

which is indep of t.



If  $L/R = CR$   
 Show that the current drawn from the source is indep of time.